

Accretion Physics

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***Black Hole Accretion:
Highly Energetic Process
Makes Relativistic Jets***



Artist's Image (credit: Lincoln Greenhill, Makoto Inoue, Jim Moran)

Review Articles

- **Black Hole Accretion: Narayan & Quataert, Sci. 307, 77-80 (2005)**
- **Hot Accretion Flows around Black Holes: Yuan & Narayan, ARAA 52, 529-588 (2014)**

Accretion Disk



- Gas with **angular momentum** goes into orbit at a large radius around the **BH**
- Gas spirals in towards the center via a sequence of near-circular orbits
- Substantial amount of **gravitational energy** is released as **heat**
- Energy comes out as **radiation, jets**
- The **EHT** makes this a timely subject

Angular Momentum & Energy

- Consider circular Keplerian orbits (to fix ideas):
- Specific angular momentum decreases inward
→ gas must lose angular momentum to spiral in

$$l = \sqrt{GM r}$$

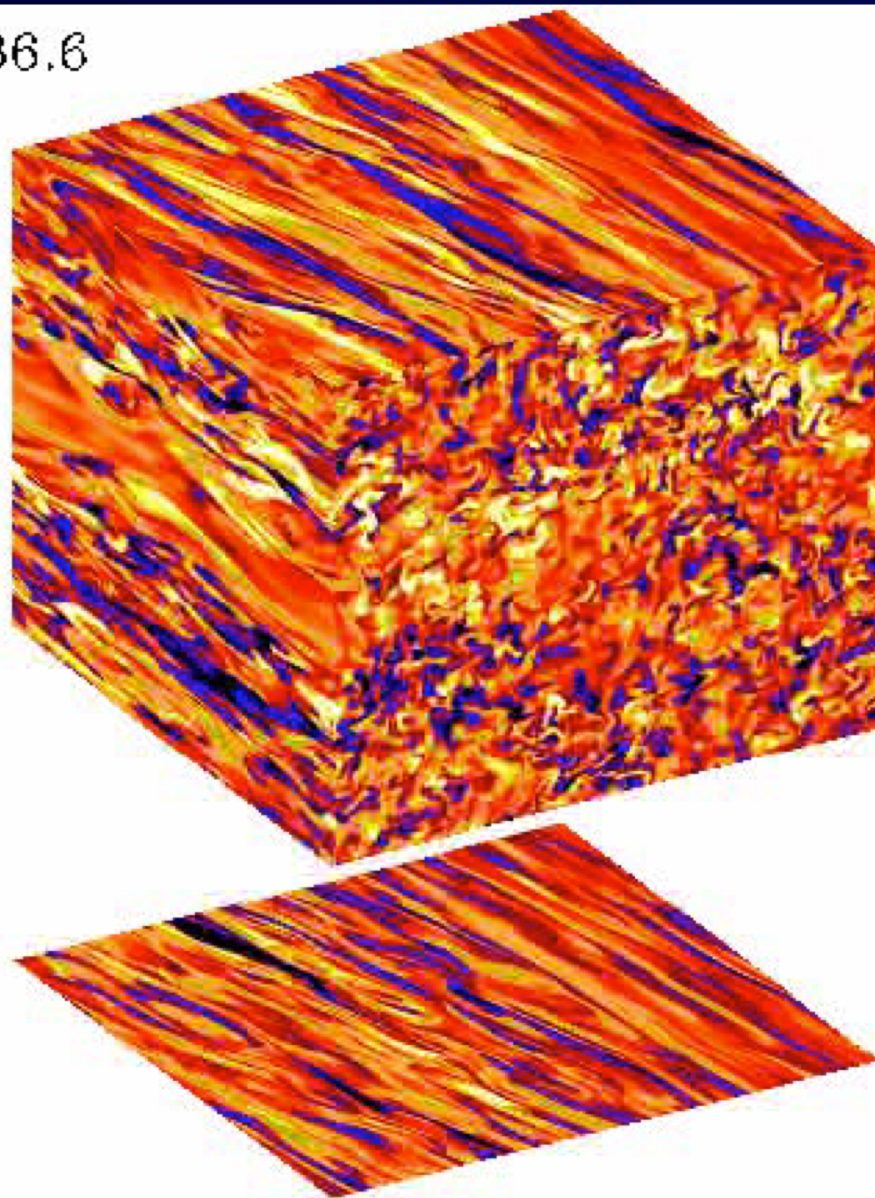
- Binding energy increases inward
→ gas must get rid of energy to remain Keplerian

$$e = GM / 2r$$

Angular Momentum Transport

- “Viscosity” in the gas can transport ang. mmtm outward, thus allowing accretion
- But microscopic viscosity is too weak
- **Major Breakthrough:** Magnetorotational instability (MRI: **Balbus & Hawley 1991**)
 - Weak magnetic field in a differentially-rotating disk causes MRI
 - Nonlinear development of MRI gives turbulence and angular mmtm transport

$t=1336.6$
 B_y



Movie Credit: Axel Brandenburg

Importance of Numerical Computations

- Numerical simulations are vital
 - Magnetic fields (MHD, MRI, and beyond)
 - Nonlinear turbulence
 - General relativity: Kerr metric
 - Webinar: Charles Gammie
- Ray-tracing in curved space-time
 - Webinar: Chi-Kwan Chan

What About Energy?

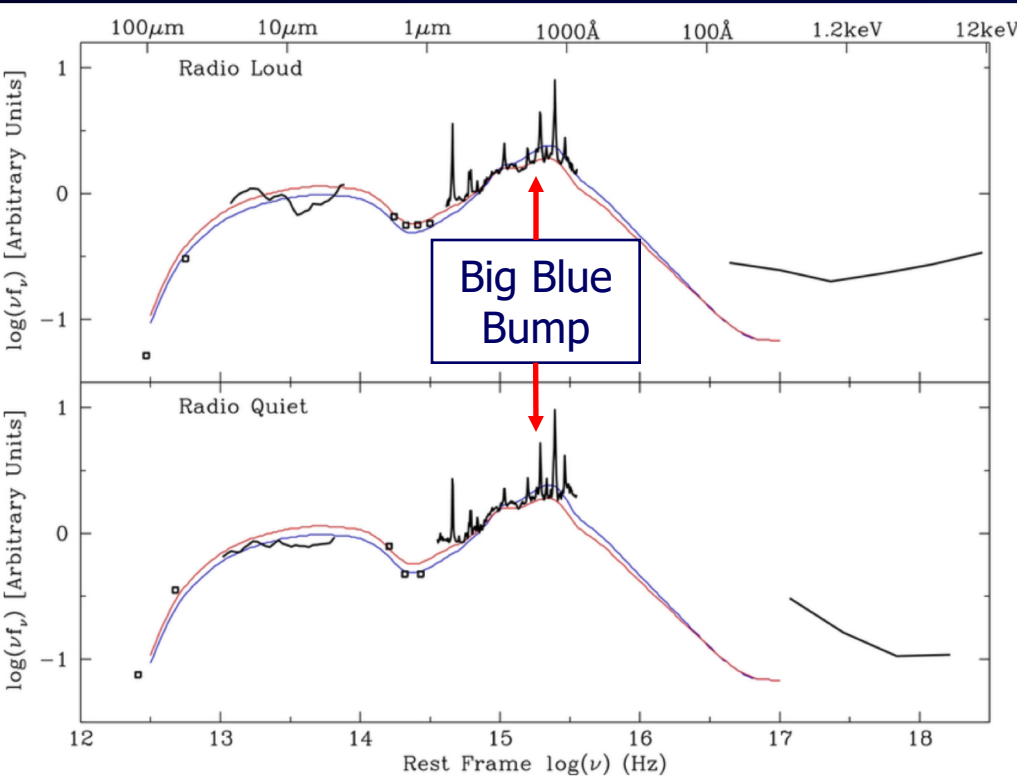
- Accretion down to the center releases a lot of energy: $\sim GM/2r$
- This energy appears as **heat** (viscous dissipation into heat is unavoidable)
- Under favorable cases, the hot gas will simply radiate away all the heat
- We then have a **Thin Accretion Disk**

Thin Accretion Disk Model

(Shakura & Sunyaev 1973; Novikov & Thorne 1973)

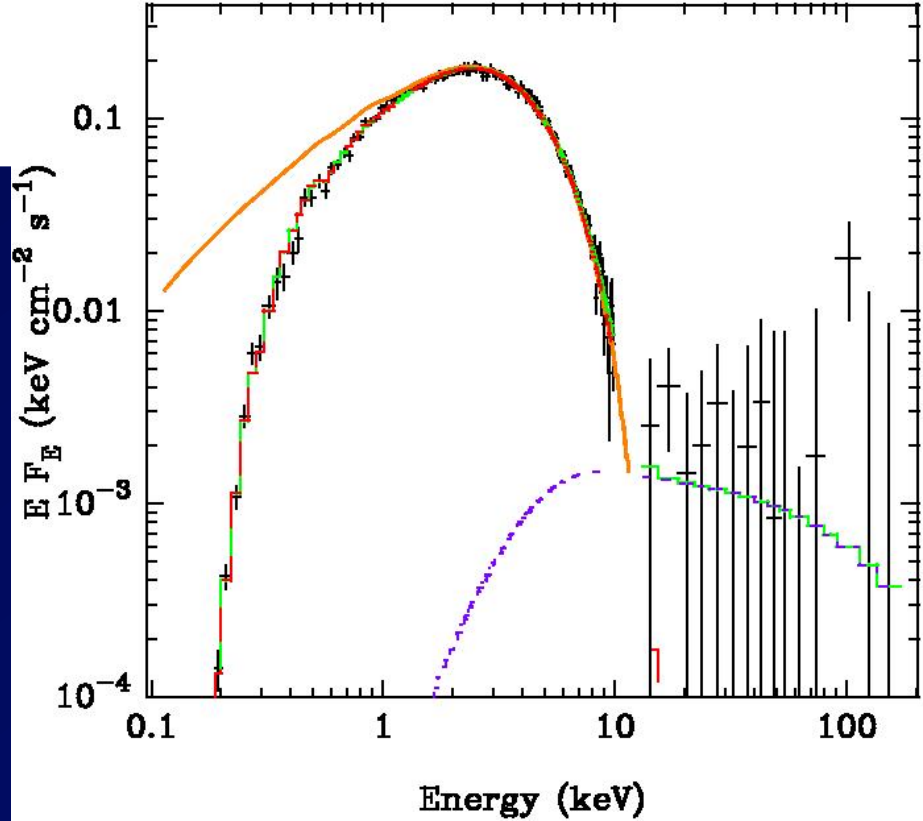


- Geometrically thin: $H \ll R$
 - Cool gas: $c_s \ll v_K$, $p \ll \rho v_K^2$ (gas pressure is negligible)
 - Keplerian rotation: $\Omega = \Omega_K$
 - Radiatively efficient: $L_{\text{acc}} \sim 0.1 \dot{M} c^2$
- Optically thick --- blackbody-like radiation
 - multi-color blackbody spectrum
- Quasars, X-ray binaries in the “soft state”



Typical spectra of luminous quasars. The Big Blue Bump in the spectrum is believed to be from a thin disk.

Thermal bump in spectrum of X-ray Binary LMC X-3 in the soft state. Modeled well with a thin disk.



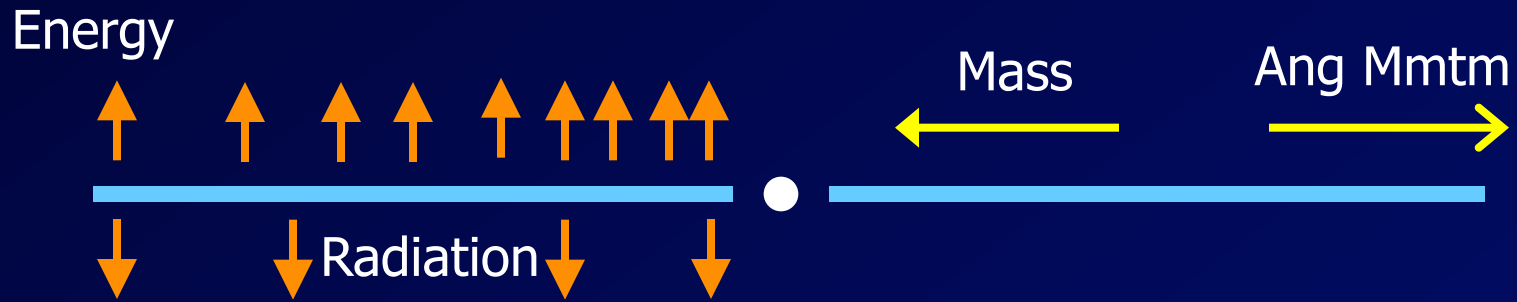
Eddington Limit

$$L_{\text{Edd}} = \frac{4\pi GMc}{\kappa} = 10^{38} \frac{M}{M_{\odot}} \text{erg s}^{-1}$$

$$\dot{M}_{\text{Edd}} = \frac{L_{\text{Edd}}}{0.1c^2} \approx 2 \times 10^{-8} \left(\frac{M}{M_{\odot}} \right) M_{\odot} \text{yr}^{-1}$$

- \dot{M}_{Edd} is estimated here assuming a fiducial radiative efficiency of **10%**

Can a Disk Exceed the Eddington Mass Accretion Rate?



- As the disk luminosity approaches L_{Edd} , radiation pressure puffs up the disk vertically
- No longer geometrically thin
 - Thin Disk Model becomes inconsistent
- Does this mean $\dot{M} > \dot{M}_{\text{Edd}}$ is impossible?
- Will the disk just blow itself apart?

Advection Saves the Day

- When $\dot{M} > \dot{M}_{\text{Edd}}$, $\sim 1L_{\text{Edd}}$ worth of luminosity is radiated from the disk surface
- Rest of the energy remains trapped in the gas and is “advectioned” with the gas
- Advection-dominated accn flow (ADAF)
- Another name: Radiatively Inefficient Accn
- Internal energy: large radiation pressure
- Geometrically thick disk
- Heavy mass loss expected

Two Solutions to Energy Problem

→ Two Kinds of Accretion

Radiatively Efficient Thin Accretion Disk

Most of the heat energy is radiated

$$L_{\text{rad}} : 0.1 \dot{M} c^2$$

Advection-Dominated ADAF Radiatively Inefficient

Most of the heat energy is retained in the gas

$$L_{\text{rad}} \ll 0.1 \dot{M} c^2$$
$$L_{\text{adv}} : 0.1 \dot{M} c^2$$

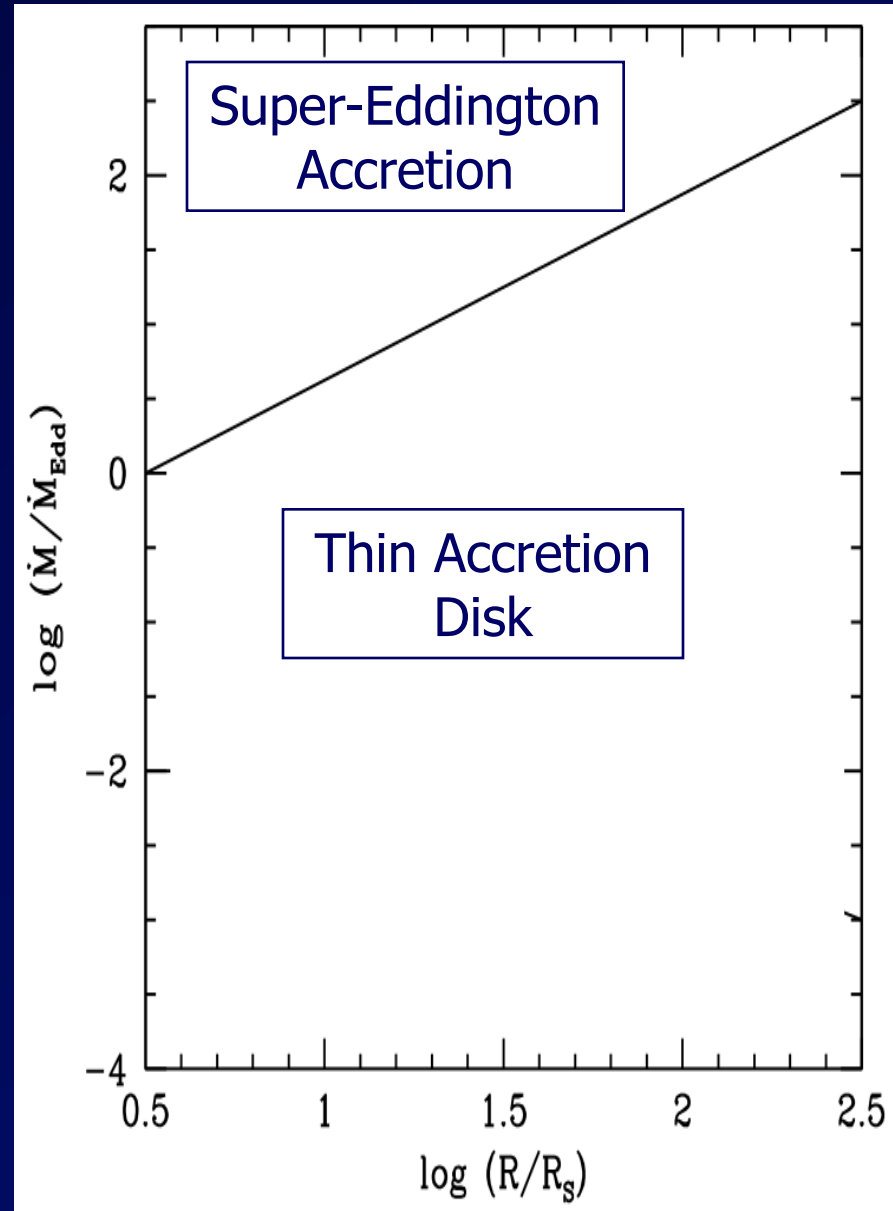
Two Distinct Accretion Regimes

Looks like a simple story:

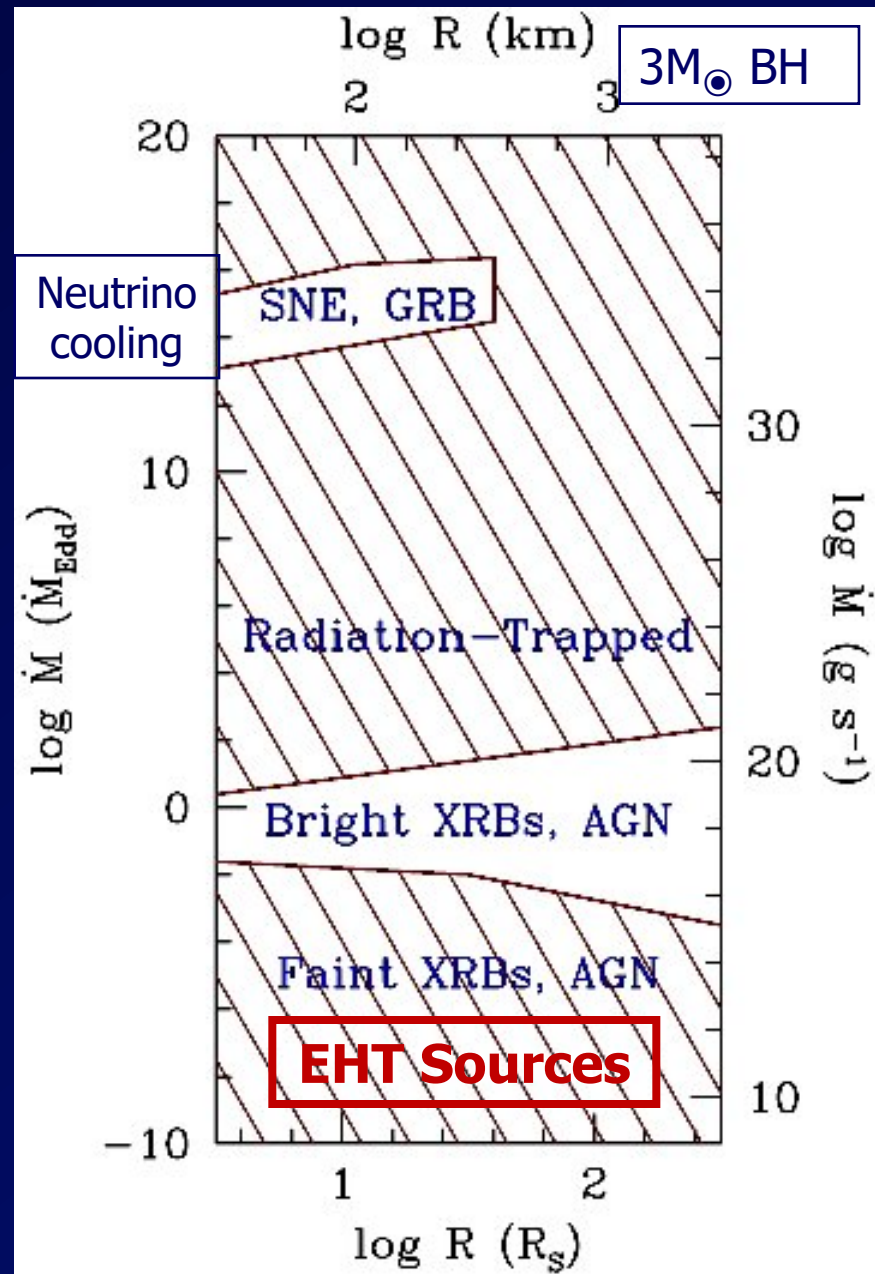
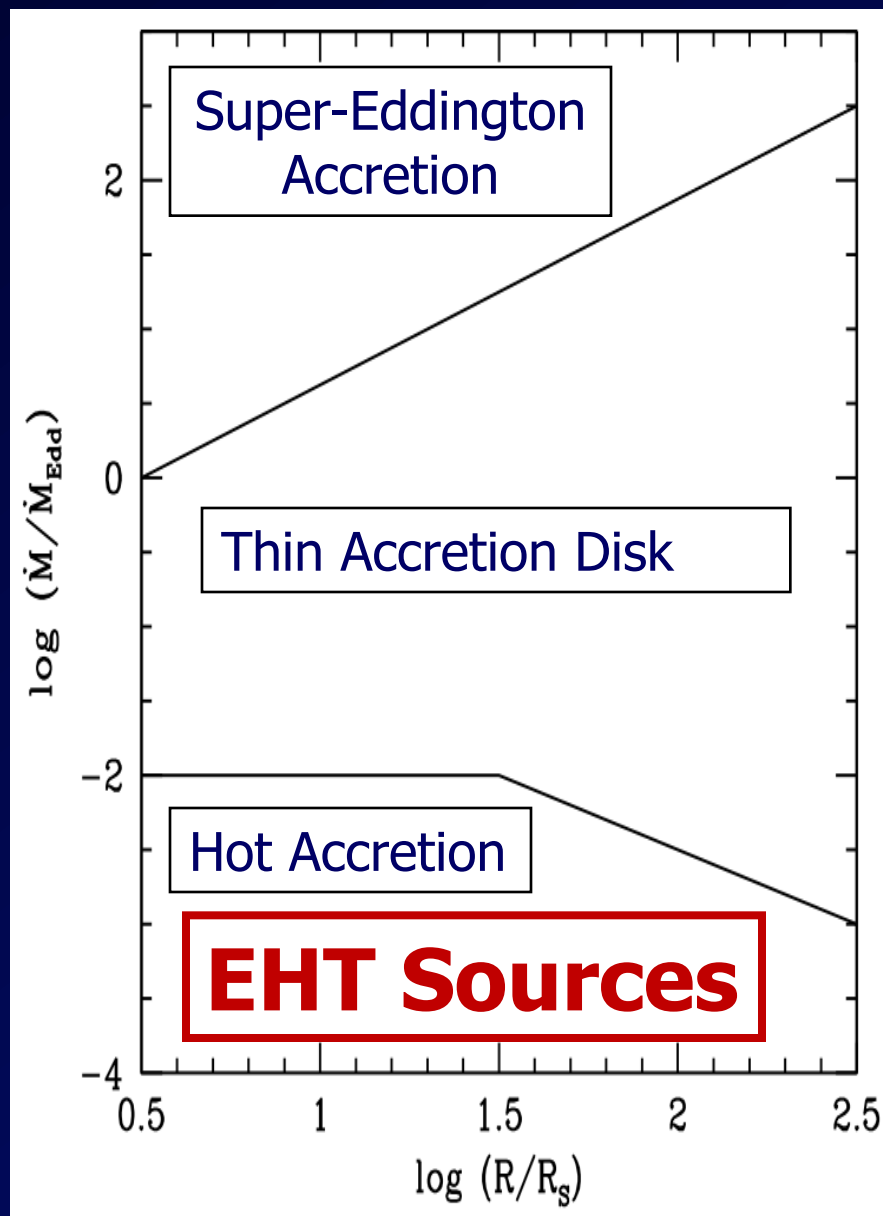
Radiatively efficient thin disk when $F_{\text{local}} < F_{\text{Edd}}$

Advection-dominated thick disk otherwise

**Actually, the story is
much more interesting!**



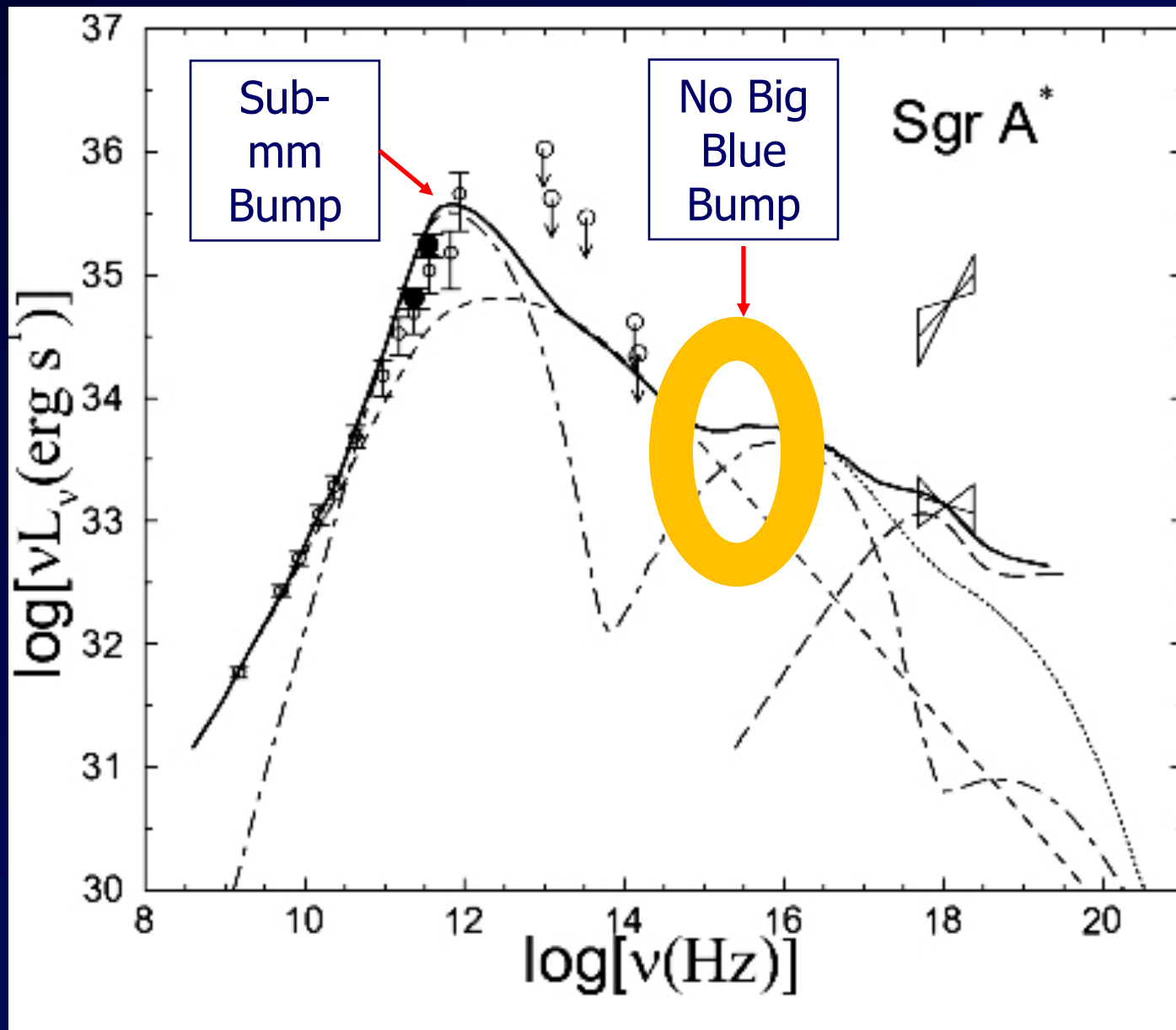
Accretion Regimes



Narayan & Quataert (2005)

A Surprise at Low \dot{M}

- For $\dot{M} \ll \dot{M}_{\text{Edd}}$, the thin disk model is perfectly good (very very thin)
- But Nature doesn't like this model!
- When $\dot{M} \lesssim 0.01 \dot{M}_{\text{Edd}}$, both supermassive and stellar-mass BHs seem to switch to a **Hot Accretion Flow**
- Thermal Big Blue Bump goes away, and a very different spectrum appears



Galactic Center BH: Sagittarius A*

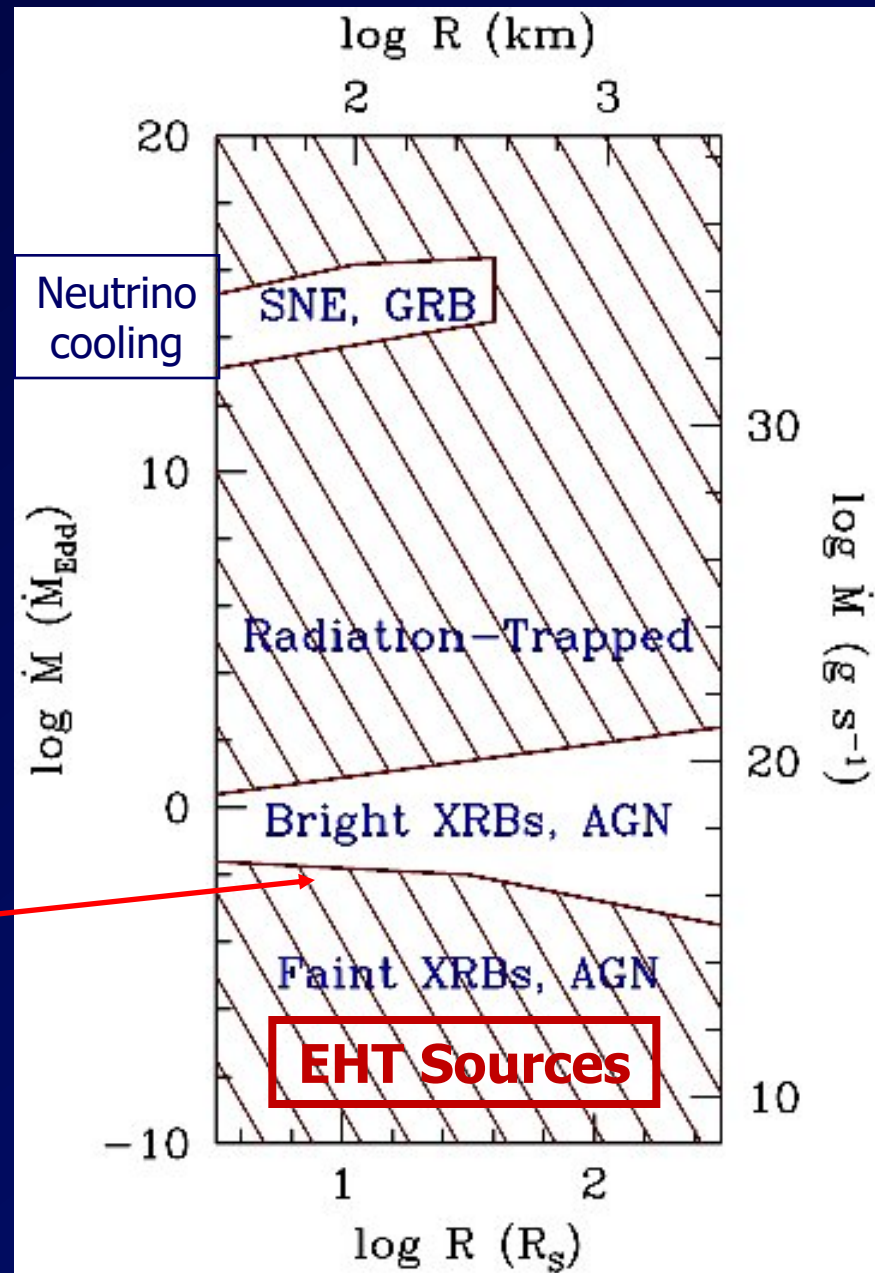
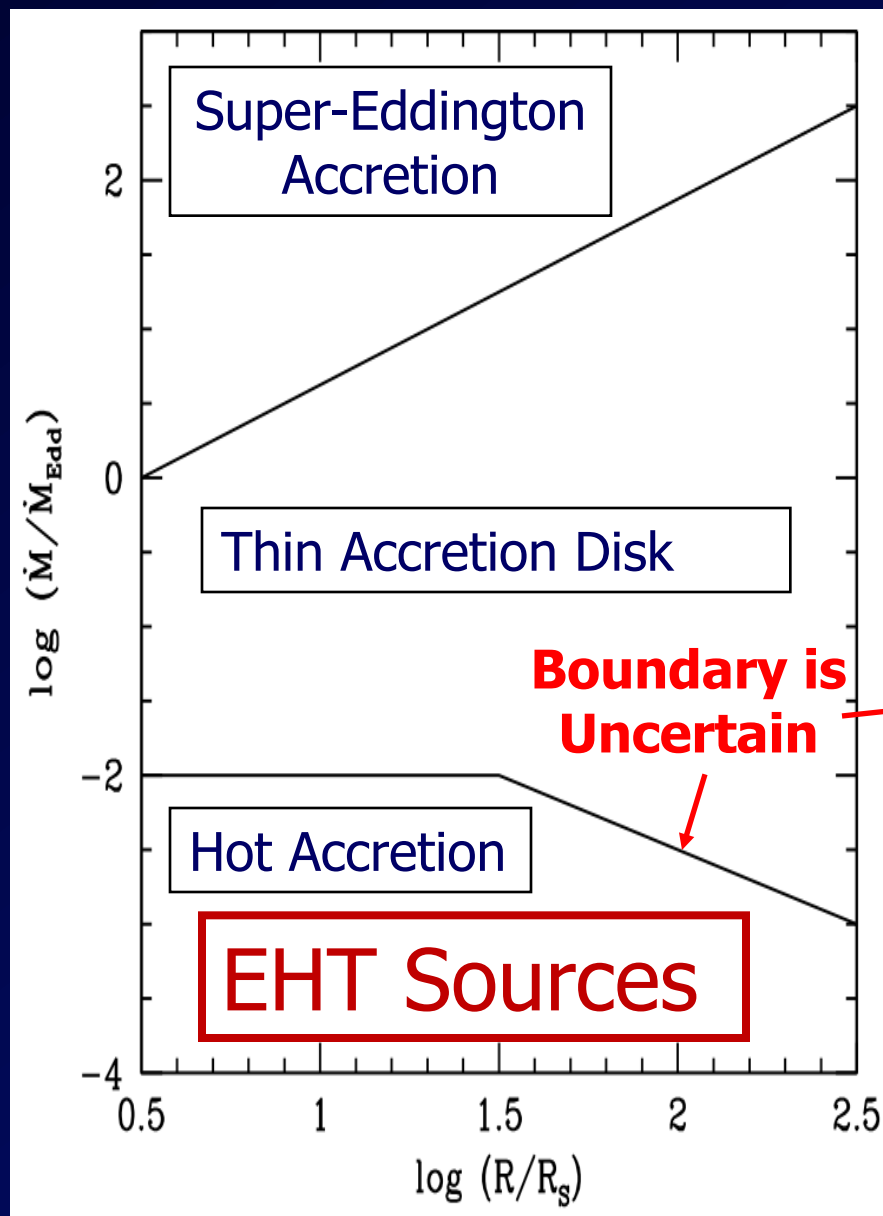
Theory: A Hot Accretion Flow Solution Appears at Low \dot{M}

- For $\dot{M} \lesssim 0.01 \dot{M}_{\text{Edd}}$, a new accretion solution appears which
 - is advection-dominated (ADAF)
 - contains very hot gas (virial temperature)
 - is geometrically very thick
 - is optically thin
- The thin disk solution is still allowed, and is perfectly viable!

So Why Does Nature Prefer Hot Accretion at Low \dot{M} ?

- Nature seems to be just waiting for the hot solution to appear and switches to it immediately
- But why is this the case?
- There is no really good answer
- It is also not clear how quickly the transition can happen

Accretion Regimes

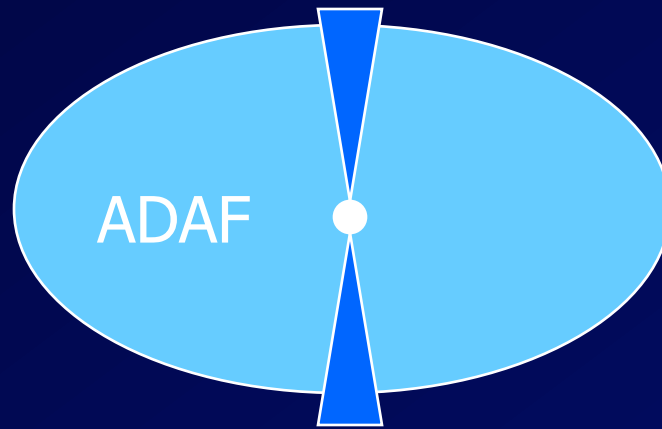


Narayan & Quataert (2005)

Hot Accretion Flow Geometry

Pure ADAF -- no thin disk at all

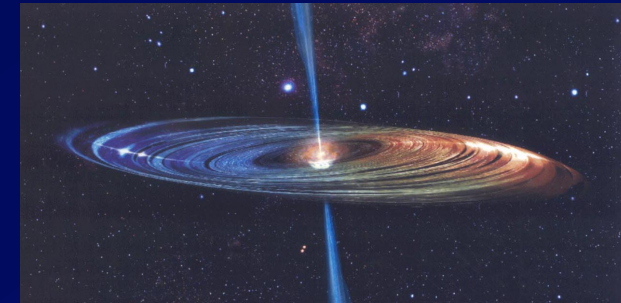
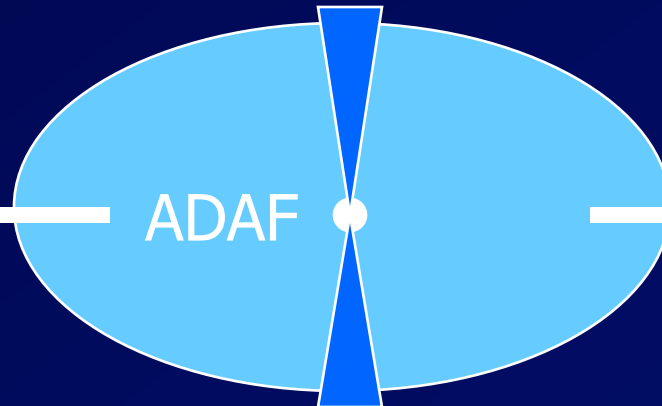
Examples: Sgr A*, perhaps M87



External Medium

Thin disk transitions to ADAF at small r

Example: NGC4258?



Thin Disk

Hot Accretion Flow Properties

- **ADAF**: Radiatively inefficient
- Energy stored in gas (gas/mag pressure)
- Gas becomes very hot: $T \sim 10^{12} \text{K}/r$
- Disk is geometrically thick: $h \sim r/2$
- Density becomes very low: $\rho \sim T^{-3/2}$
- Optically thin: $\tau_{\text{es}} \ll 1$ (we can see the BH!)
- Cooling rate is low: even though T is large, ρ is very small (brems: $\rho T^{1/2}$)
- Low cooling \rightarrow radiatively inefficient

Two-Temperature Plasma

- At the very low ρ found in a hot accretion flow, the plasma is collisionless
- Likely to become two-temperature with $T_e \ll T_i$ (e.g., $T_i \rightarrow 10^{12}\text{K}$, $T_e < 10^{11}\text{K}$)
- This enhances radiative inefficiency
- Can the EHT “prove” that the plasma in M87, Sgr A* is two-temperature?

Radiation from a Hot ADAF

- Hot electrons with temperature $>10^{10}\text{K}$ radiate primarily via
 - Synchrotron
 - Bremsstrahlung
 - Comptonization:
 - Synchrotron self-Compton
 - External soft photons (if there is an outer disk)
- Ions ($>10^{11}\text{K}$) hardly radiate (pions?)

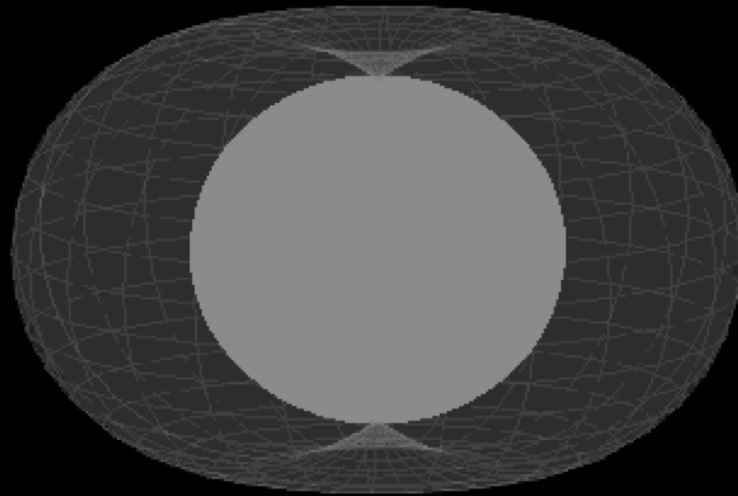
Energy Problem: To Accrete or Not?

- Without radiative cooling, accretion energy is stored as pressure (gas, B)
- Gas is only weakly bound to the BH: pressure balances gravity
- Two limits of energy advection:
 - All the gas accretes into the BH
 - (Almost) all the gas is ejected
- The truth is probably in between
- Waiting for simulations to give the answer

ADAFs and Relativistic Jets

- Numerical simulations show that hot accretion flows produce **powerful relativistic jets**
- Physics consistent with **Blandford & Znajek (1977)**: **B** plus **BH spin**

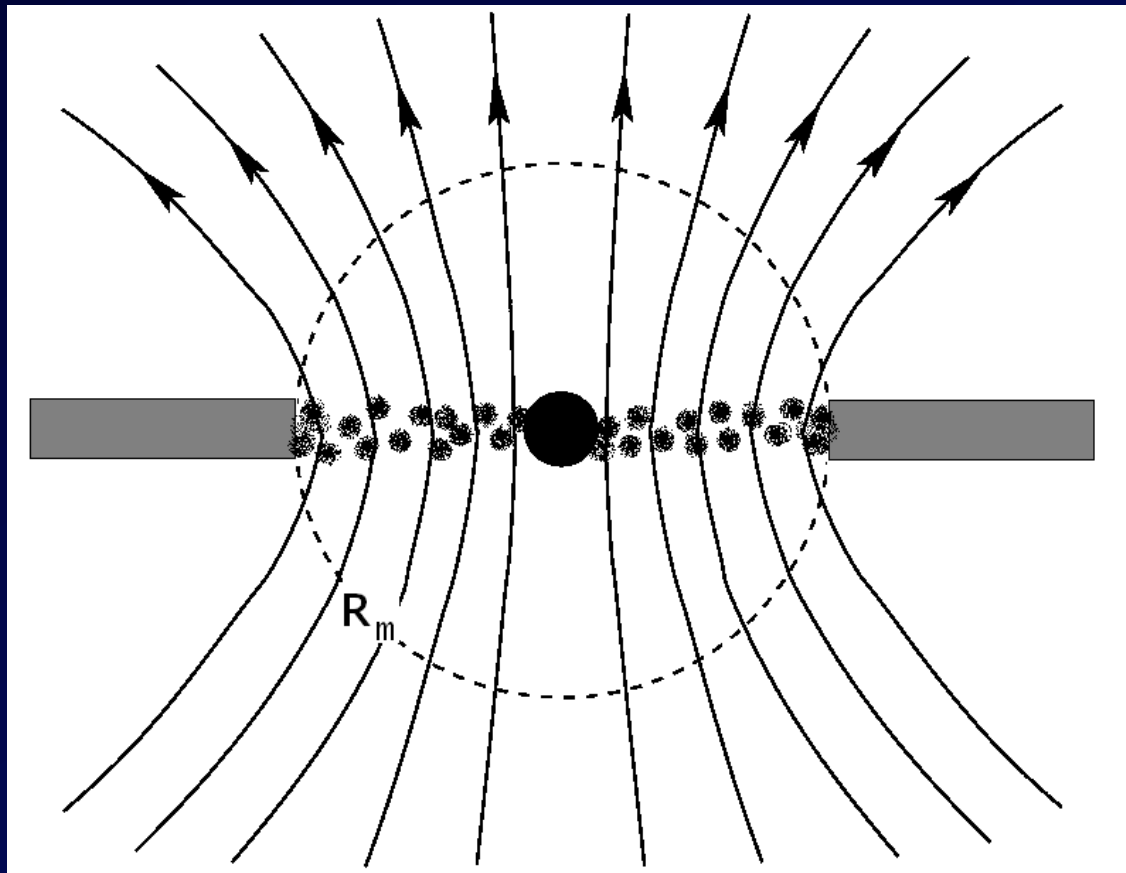
$$P_{\text{jet}} \approx \Phi_{\text{BH}}^2 a_*^2$$



Semenov et al. (2004)

Black Hole Spin Energy!

- $P_{\text{jet}} = 1.4 \dot{M} c^2$ for a simulation with $a_* = 0.99$ (Tchekhovskoy 2011)! How?!
- Jet derives most (all?) of its power from the spin energy of the BH
- Can we have arbitrarily large jet power ($100 \dot{M} c^2$) by simply increasing B ?
- No! There is a limiting field strength: Magnetically Arrested Disk (MAD)



MAD: The field is as strong as it can be and yet permit gas to accrete. Accretion with weaker fields is called **SANE**

$$\phi_{\text{MAD}} = \Phi_{\text{BH}} / \sqrt{\dot{M} r_g^2 c} \approx 50$$

EHT and Jets

- Can observations with the EHT “prove” that the jet in M87 is powered by the Blandford & Znajek mechanism?
- Can observations with the EHT “explain” why Sgr A* apparently has no jet (or at best a very weak one)?

Hot Accretion Thermodynamics is Highly Uncertain

- Plasma in a Hot ADAF is collisionless
 - Electrons/protons do their own thing
- Non-equilibrium: Two-temperature + ...
- Each particle remembers its heating history and radiates accordingly
- We need to understand plasma heating processes to interpret observations
- Both a curse and an opportunity

Collisionless
Accretion flow →
Viscous dissipation
→ Heat

Ions

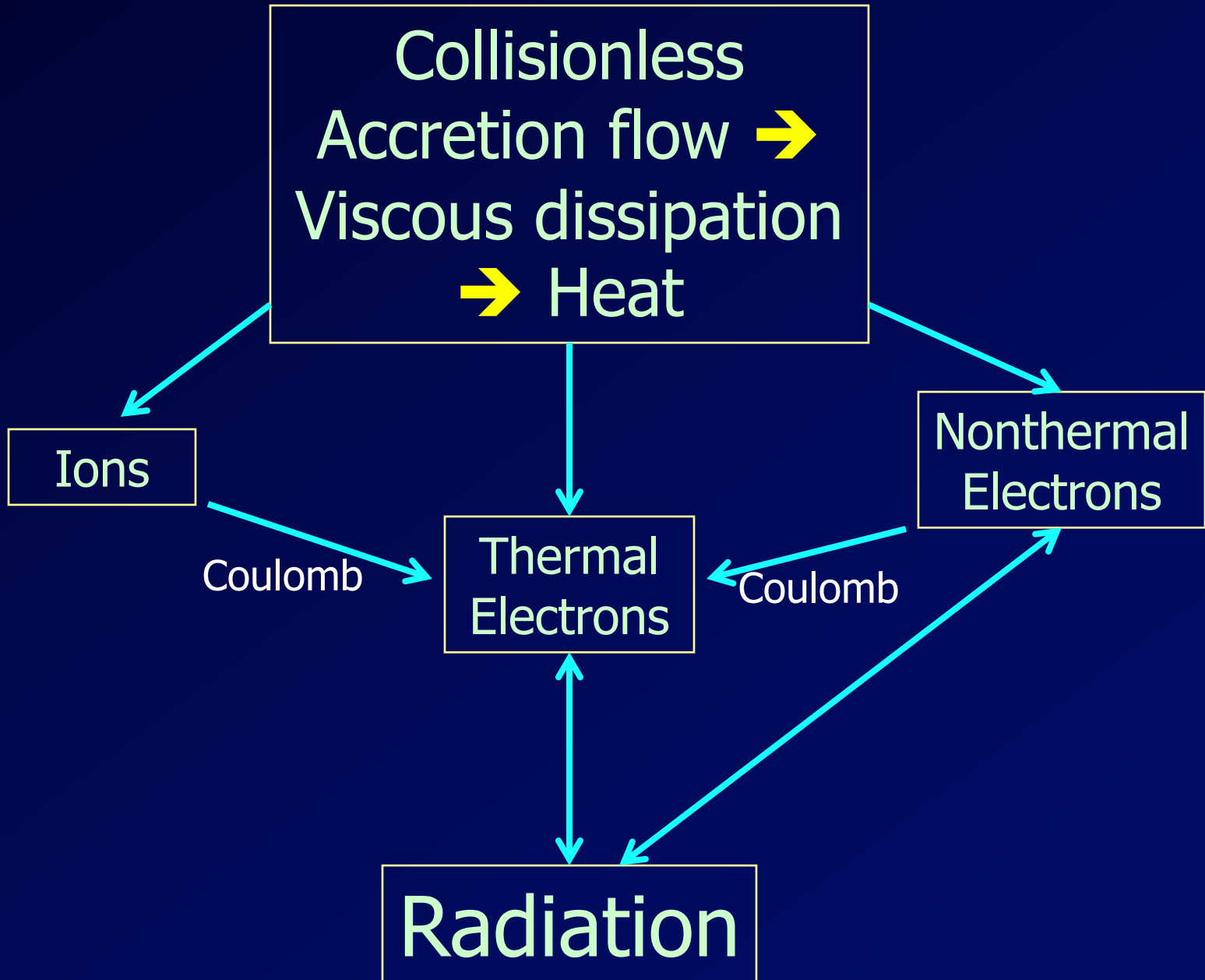
Nonthermal
Electrons

Coulomb

Thermal
Electrons

Coulomb

Radiation



Particle Heating/Acceleration

- Multiple heating/acceleration channels
 - Compression, shear, instabilities
 - Turbulent cascade and dissipation
 - Magnetic reconnection
 - Shocks

Upcoming Webinars

- **Gammie:** GRMHD Numerical Simulations
- **Chan:** Ray Tracing
- **Quataert:** Subgrid/Plasma Physics
- **Berti:** Kerr Black Hole and Beyond

Concluding Comments

- This webinar gave a broad survey of
 - Accretion Physics, esp. Hot Accretion Flows
 - Connection to Relativistic Jets
- The EHT has triggered renewed interest
 - EHT motivates re-examining our ideas
 - EHT can test our understanding
 - EHT can/will come up with surprises

***Please Complete
the Survey!***

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